Magnetism and Spintronics: 1

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#### **Definitions**

#### B:

- magnetic induction (usually means by magnetic field).
- most physically significant quantity what shows up in Lorentz force law, in determining NMR frequencies, etc.
- Unit is the Tesla. Earth's magnetic field = 6 x 10<sup>-5</sup> T.
- More convenient cgs unit is the Gauss = 10<sup>-4</sup> T.
- Important boundary condition:  $\nabla \cdot \mathbf{B} = 0$

### **Definitions in SI**

### **H**:

- the magnetic field.
- Caused by currents of free charge.
- Unit is the Amp/m.
- Important relation:  $\nabla \times \mathbf{H} = \mathbf{J}$
- With no magnetic materials around, B=μ<sub>0</sub>H

#### M:

- the magnetization.
- magnetic moment per unit volume of a material.
- Unit is the Amp/m.
- For a material with a permanent magnetization  $M_0$ ,

"permeability of free

space" =  $4\pi \times 10^{-7} \text{ Tm/A}$ 

$$\mathbf{M} = \chi \mathbf{H} + \mathbf{M}_0$$

1

magnetic susceptibility

# Susceptibility and permeability

Combining effects of external currents and material response,

$$\mathbf{B} = \mu_0(\mathbf{H} + \mathbf{M})$$

We define the *permeability* by:  $\mathbf{B} = \mu \mathbf{H}$ 

So, for a material with a magnetic susceptibility  $\chi$ ,

**B**=
$$\mu$$
**H**= $\mu_0$ (1+ $\chi$ )**H**  $\rightarrow \mu$ = $\mu_0$ (1+ $\chi$ )

Note that for real materials  $\chi$  and  $\mu$  are tensorial. Relative permeability is defined as  $\mu_r = \mu/\mu_0$ .  $\longrightarrow$   $\mathbf{B} = \mu_0 \mu_r \mathbf{H}$ 

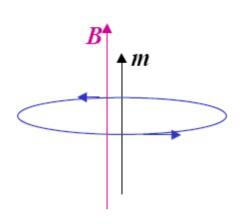
Susceptibility is most useful when discussing *diamagnetic* ( $\chi$  < 0) and *paramagnetic* ( $\chi$  > 0) materials, rather than systems with nonzero  $M_0$ .

# Diamagnetism

Some materials develop a magnetization that is *antialigned* with the applied external field H.

Such materials are *diamagnetic*, and have  $\chi < 0$ .

Simple *classical* picture for diamagnetism: Lenz's Law



Try ramping up  $\mathbf{B} = \mu_0 \mathbf{H}$ .

Result is a circumferential electric field that *opposes* the direction of the current in the loop.

This would act to reduce the dipole moment along *H*, and would be diamagnetic.

Correct quantum treatment involves 2nd order perturbation theory - can end up with either sign, depending on particulars of atoms. *Larmor* diamagnetism or *Van Vleck* paramagnetism.

# Paramagnetism

Also common is paramagnetism, when  $\chi > 0$ .

Two common origins of paramagnetism:

- Curie paramagnetism localized moments free to flip.
- Pauli paramagnetism requires "free" electrons in a metal.

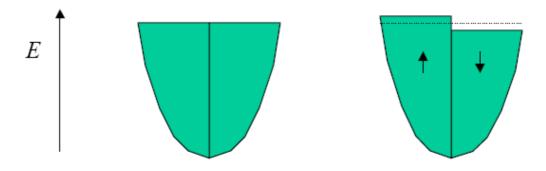
### Curie paramagnetism

Alignment of spin with external field lowers spin energy.

Calculation from statistical physics.

$$\longrightarrow \chi \propto 1/T$$

# Pauli paramagnetism



Starting from unpolarized electrons, applying a field **B** shifts the Fermi level for spin-up and spin-down electrons oppositely!

$$N_{up} = V \int d\varepsilon \frac{v(\varepsilon)}{2} f(\varepsilon + \mu_{\scriptscriptstyle B} B) \qquad N_{\scriptscriptstyle down} = V \int d\varepsilon \frac{v(\varepsilon)}{2} f(\varepsilon - \mu_{\scriptscriptstyle B} B)$$

Taylor expanding to find the difference,

$$M = \frac{\mu_B}{V} (N_{up} - N_{down}) \approx \mu_B^2 \cdot v(E_F) \mu_0 H$$

Only good for metals - can't have gap at Fermi surface.

# Ferromagnetism - toy model

Start from Curie paramagnetism picture - local moments (spins), but now allow them to respond to the *local* magnetic field at their position.

Assume 
$$M = \chi_0 (H + H_m)$$
 "molecular field", =  $\eta M$ 

$$M(1-\eta\chi_0)=\chi_0H$$

At high temperatures, 
$$\chi_{eff} = \frac{\chi_0}{(1 - \eta \chi_0)}$$

Recall that  $\chi_0 \sim 1/T$ , so we find at high temperatures

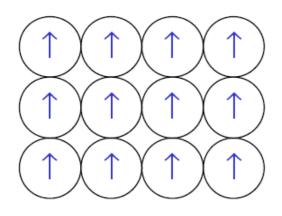
$$\chi_{eff} \sim \frac{1}{T - T_c}$$
 Curie-Weiss law

At  $T_c$ , the Curie temperature, the susceptibility diverges! Spontaneous magnetization = ferromagnetism.

### Ferromagnetism, microscopic picture

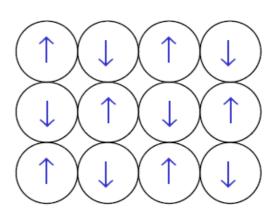
Exchange energy caused by interaction of charged Fermions

$$E_{ex} = -J\mathbf{S}_1 \cdot \mathbf{S}_2$$



Ferromagnetism J>0

$$\psi(1,2) = \varphi(1,2)S(1,2)$$
orbital spin

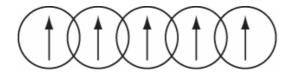


Antiferromagnetism J<0

S parallel -> φ antiparallel, lowers Coulomb enenergy

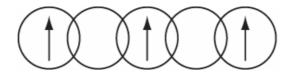
## Exchange interactions

Direct exchange



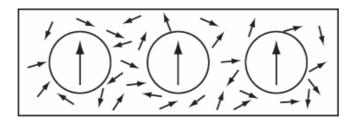
charge distribution of magnetic ions overlap

Super-exchange



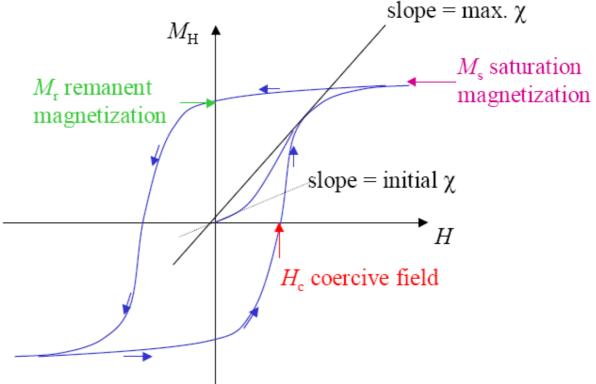
Magnetic ions interact by charge overlap with same non-magnetic ions

Indirect exchange



Magnetic ions interaction mediated by interaction with conduction elections. RKKY interaction

# Ferromagnetism: M vs H hysteresis



"Hard" ferromagnetic materials: large values of  $H_c$ , large  $M_r$ .

"Soft" ferromagnetic materials: small values of  $H_c$ , small  $M_r$ .

### Summary

- Most insulators are weakly diamagnetic.
- Metals can be either.
- Ferromagnetism caused by exchange effects.
- Susceptibilities are usually quoted *per molar volume* rather than in their dimensionless form, for experimental reasons.
- Strictly speaking, susceptibilities and permeabilities are defined as derivatives.